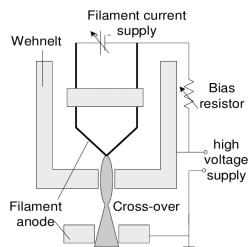
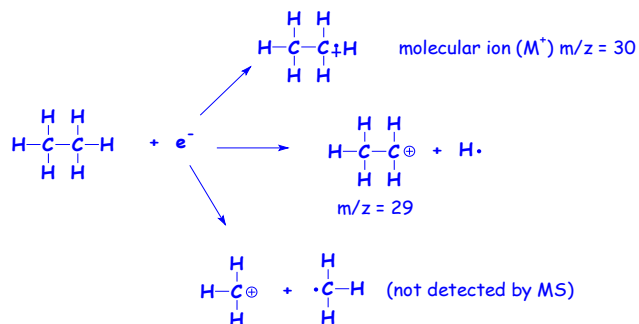


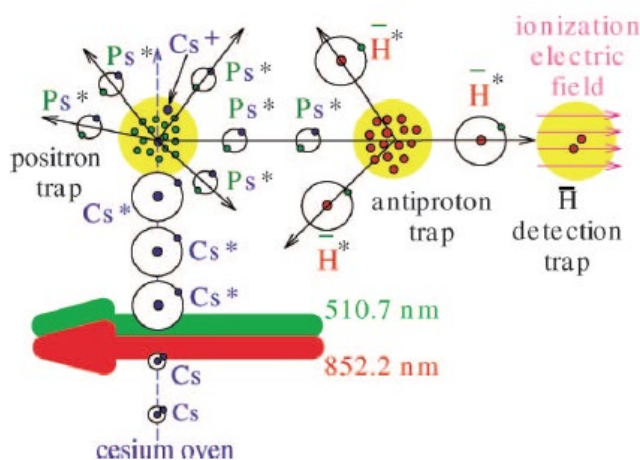
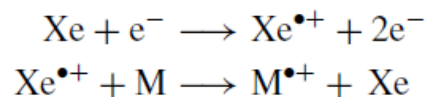
Mass spectrometry - ionization methods

1. Electron ionization

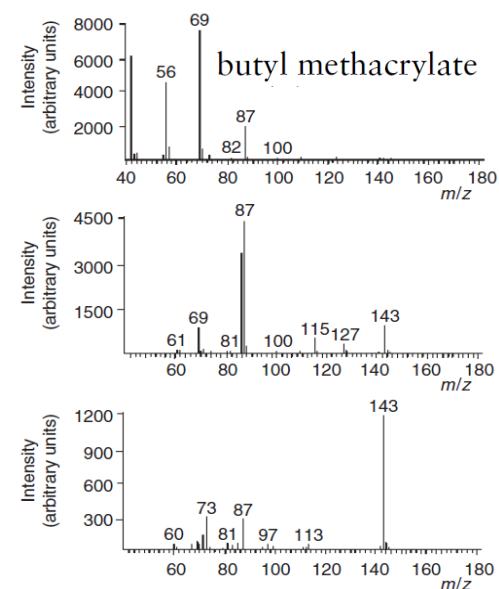


2. Chemical ionization:

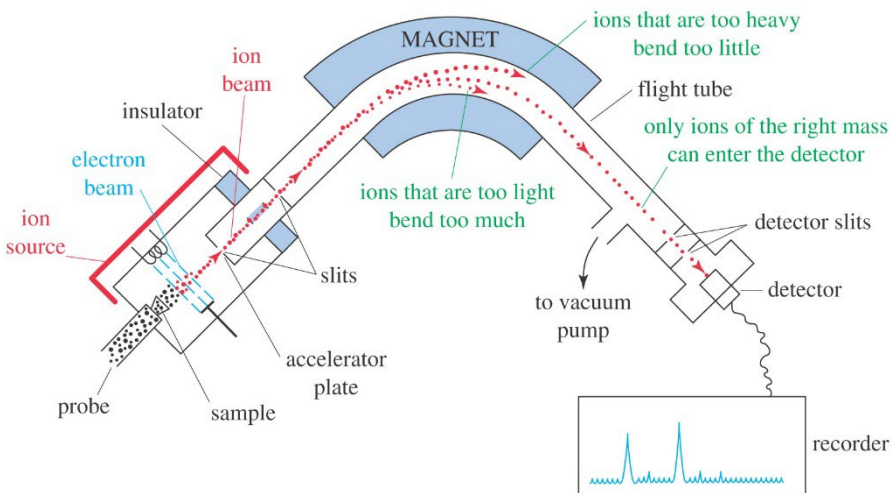
a) Charge transfer



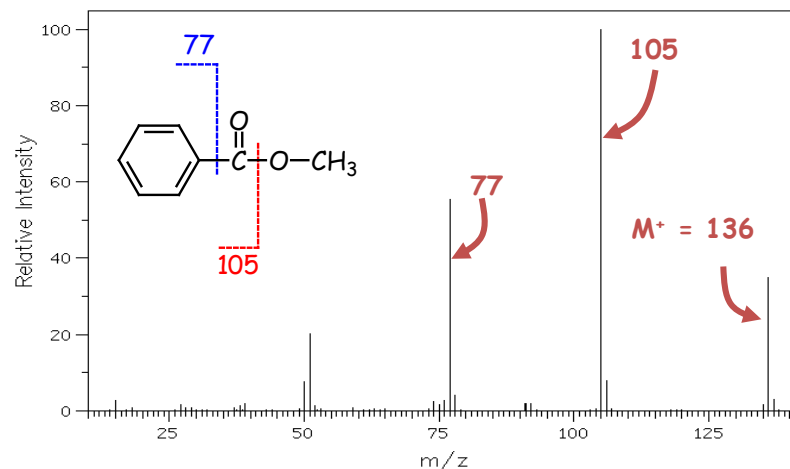
b) proton transfer



Mass spectrometry



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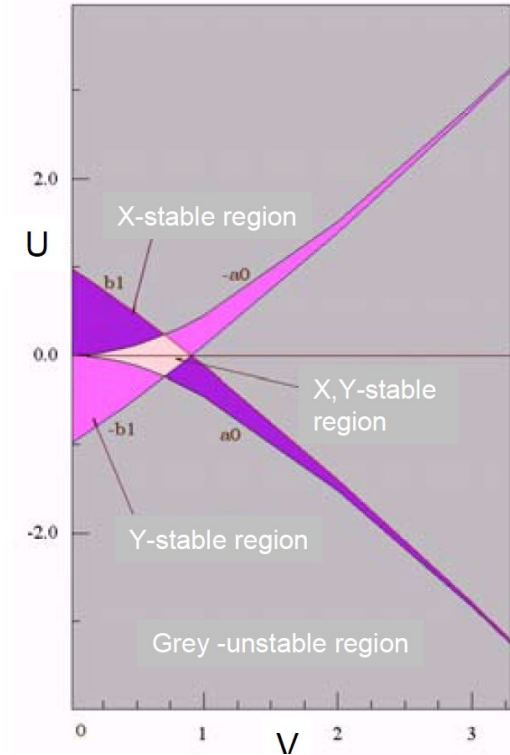
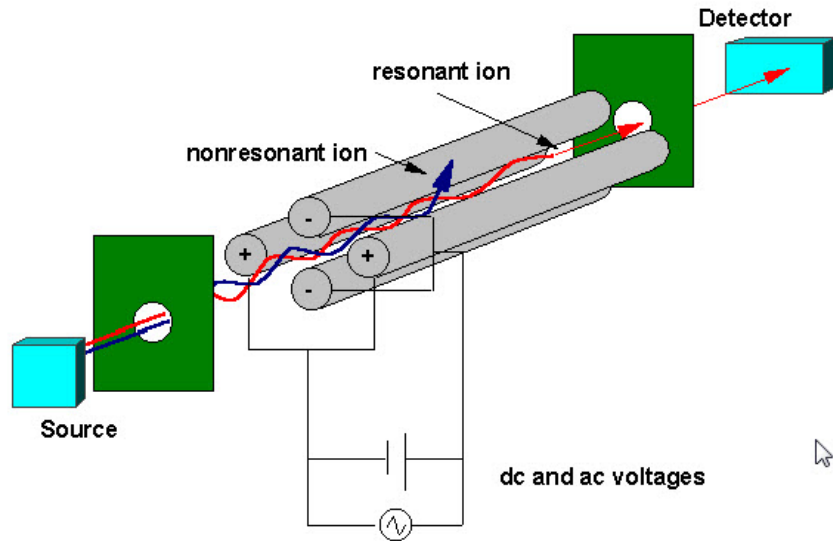


$$qvB = \frac{mv^2}{r} \quad \text{or} \quad mv = qBr \quad \frac{m}{q} = \frac{r^2 B^2}{2V_s}$$

Most elements occur naturally as a mixture of isotopes.

- The presence of significant amounts of heavier isotopes leads to small peaks that have masses that are higher than the parent ion peak.
- M+1 = a peak that is one mass unit higher than M⁺
- M+2 = a peak that is two mass units higher than M⁺

Quadrupole Mass Filter



Mathieu's equations

$$\frac{d^2x}{dt^2} = -\left(\frac{e}{m}\right) \frac{[U + V \cos(\omega t)]}{r_0^2} x,$$

$$\frac{d^2y}{dt^2} = \left(\frac{e}{m}\right) \frac{[U + V \cos(\omega t)]}{r_0^2} y,$$

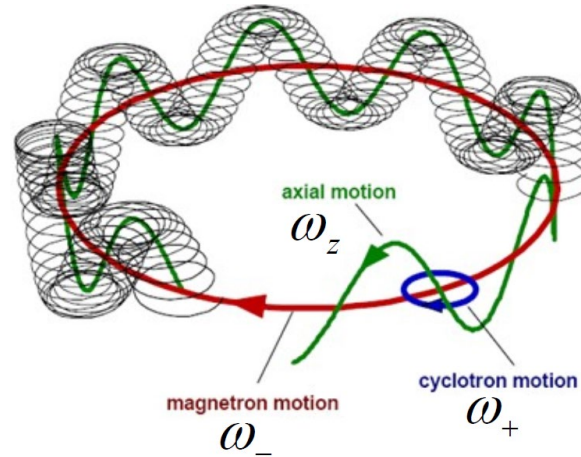
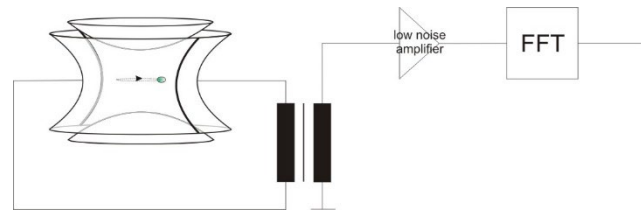
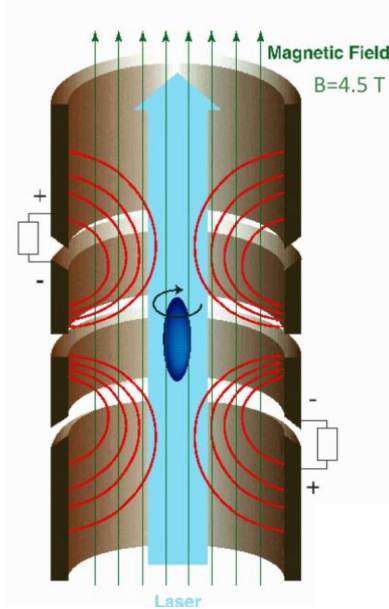
$$\frac{d^2z}{dt^2} = 0.$$

$$\frac{d^2u}{d\xi^2} + (a_u - 2q_u \cos 2\xi) u = 0$$

$$a_u = a_x = -a_y = \frac{8zeU}{m\omega^2 r_0^2}$$

$$q_u = q_x = -q_y = \frac{4zeV}{m\omega^2 r_0^2}$$

Penning trap mass spectrometry



$$\omega_c = \frac{q}{m} B$$

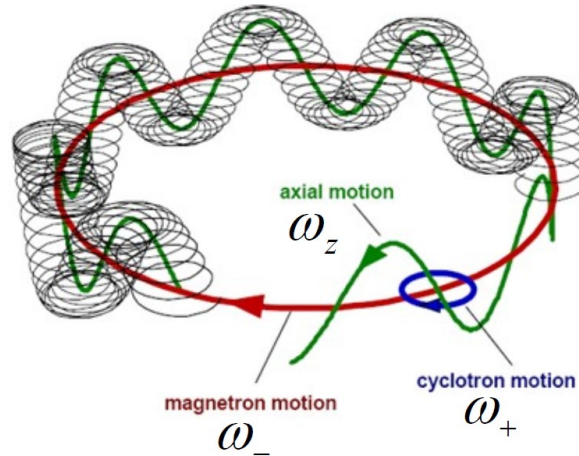
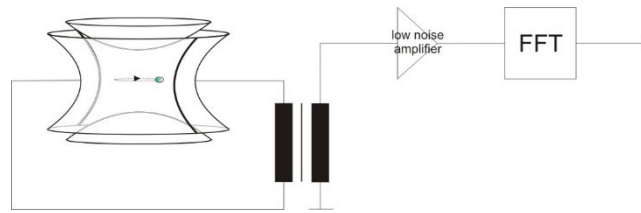
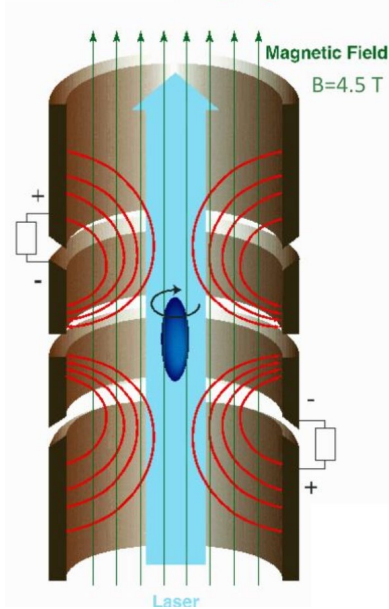
$$\omega_z = \sqrt{\frac{qU_{dc}}{md^2}},$$

$$\omega_+ = \frac{\omega_c}{2} + \sqrt{\frac{\omega_c^2}{4} - \frac{\omega_z^2}{2}}$$

$$\omega_- = \frac{\omega_c}{2} - \sqrt{\frac{\omega_c^2}{4} - \frac{\omega_z^2}{2}}$$

$$\frac{|q|}{m} B^2 > 2 \frac{|U_{dc}|}{d^2} \text{ and } qU_{dc} > 0$$

Penning trap mass spectrometry



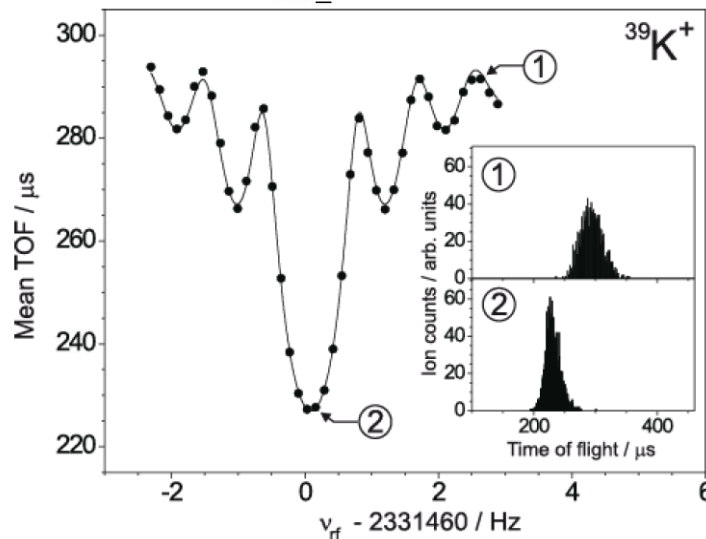
$$\omega_c = \frac{q}{m} B$$

$$\omega_z = \sqrt{\frac{qU_{dc}}{md^2}}$$

$$\omega_+ = \frac{\omega_c}{2} + \sqrt{\frac{\omega_c^2}{4} - \frac{\omega_z^2}{2}}$$

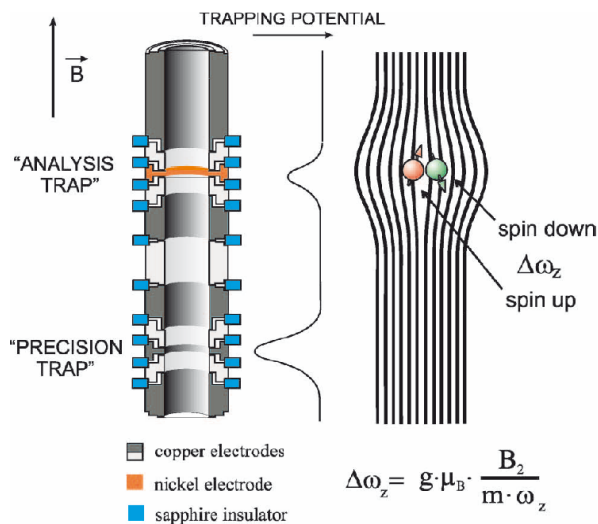
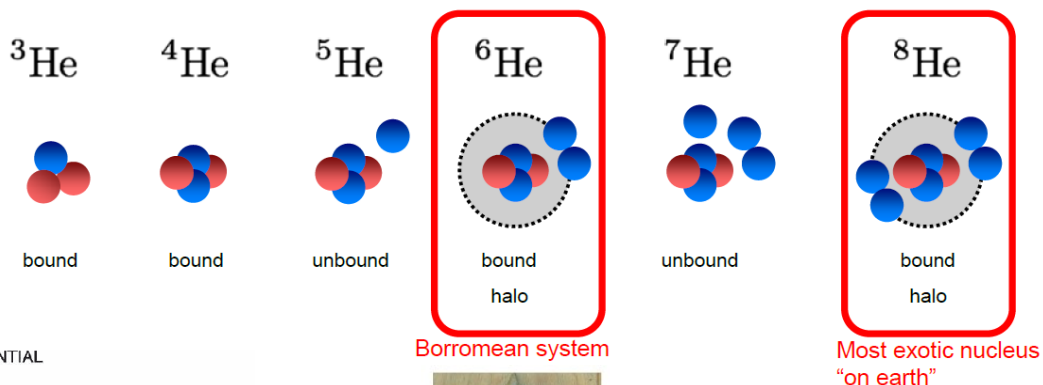
$$\omega_- = \frac{\omega_c}{2} - \sqrt{\frac{\omega_c^2}{4} - \frac{\omega_z^2}{2}}$$

$$\frac{|q|}{m} B^2 > 2 \frac{|U_{dc}|}{d^2} \text{ and } qU_{dc} > 0$$

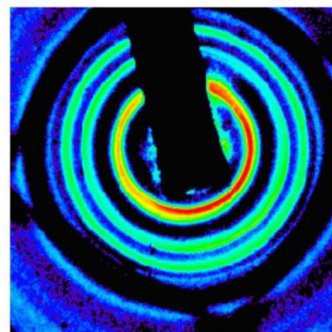


Penning trap mass spectrometry

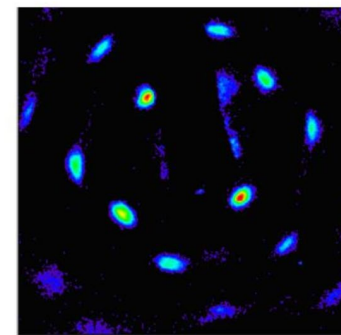
Binding energy measurements



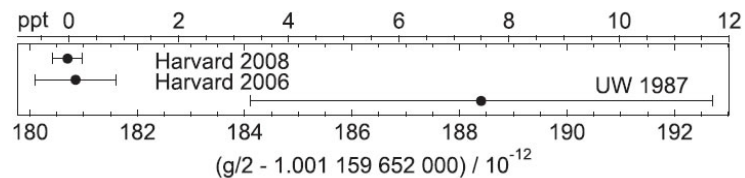
time averaged Bragg scattering



camera strobed by the rotating wall



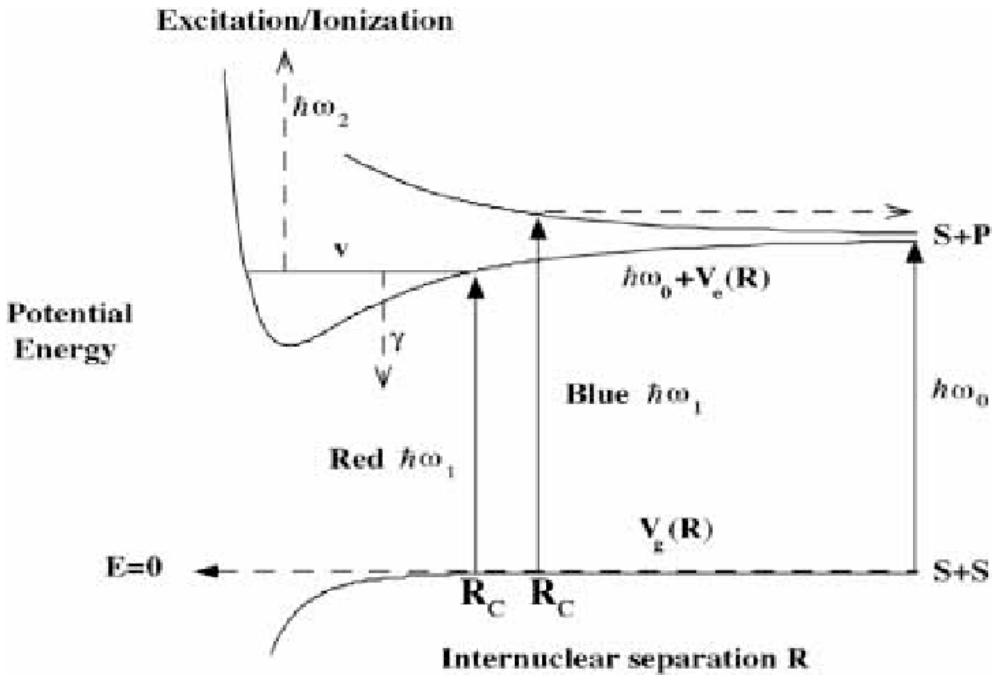
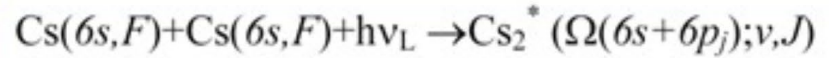
- determine if crystal pattern due to 1 or multiple crystals
- enables real space imaging of ion crystals



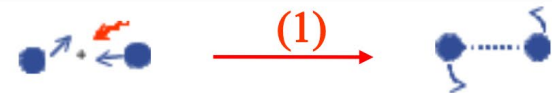
$$g/2 = 1.001\ 159\ 652\ 180\ 73(28) \quad [0.28 \text{ ppt}].$$

Electron/proton g-factor measurements

Photoassociation spectroscopy

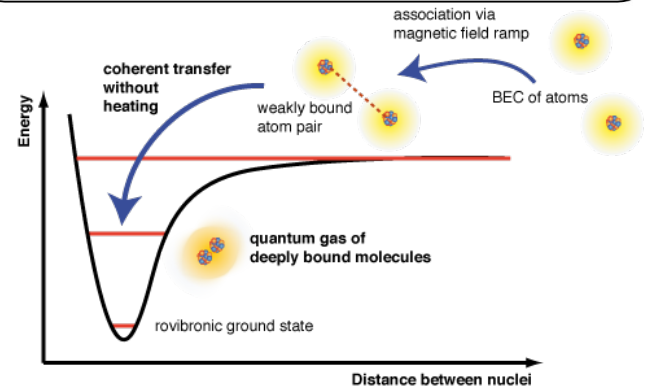
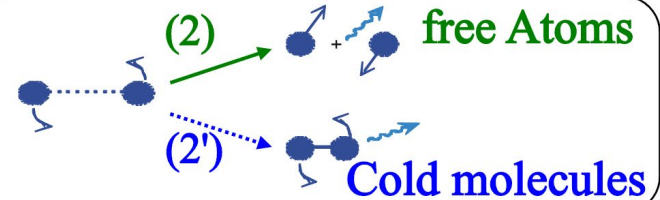


Photoassociation



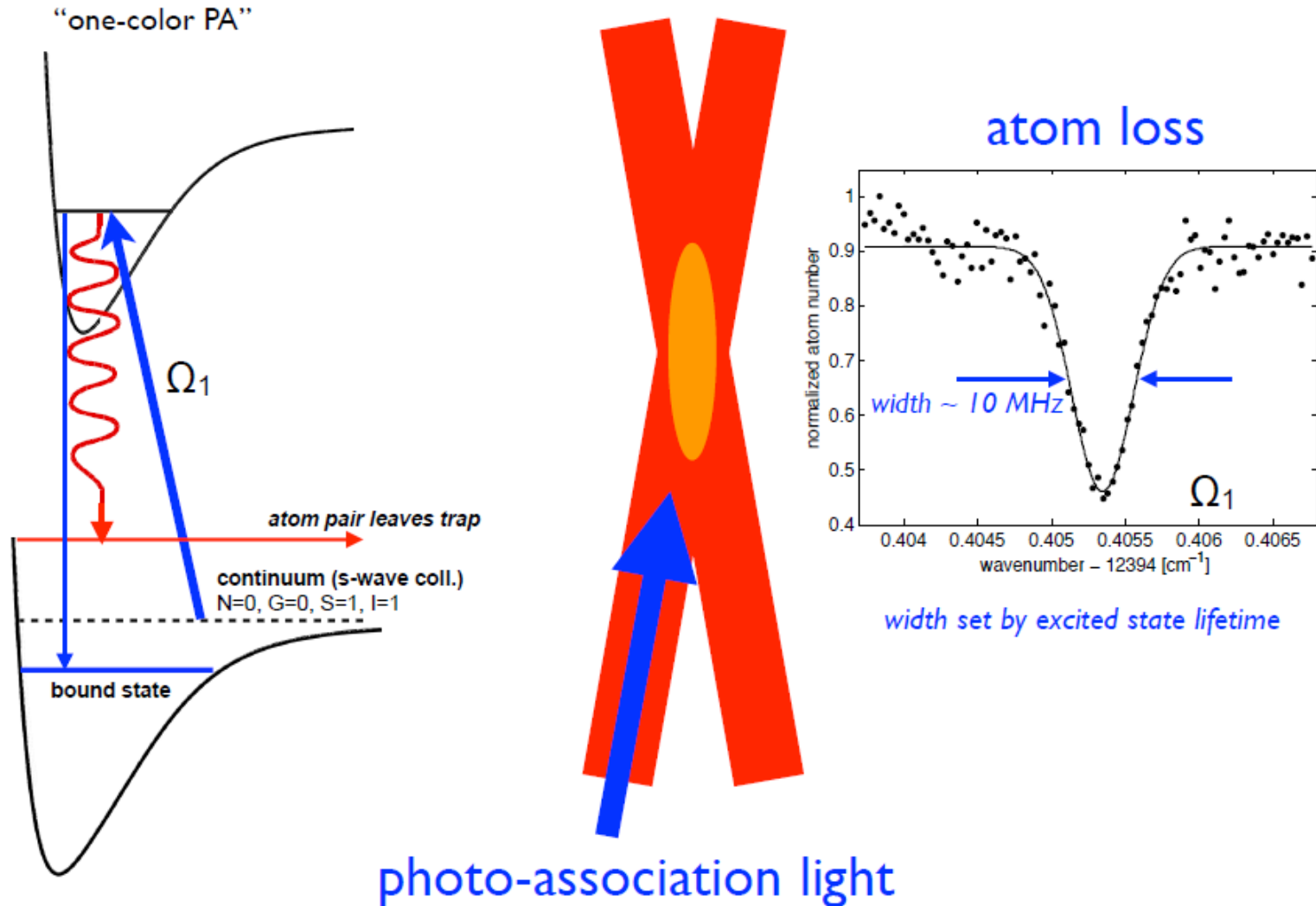
Resonant photon absorption
by two cold atoms ($T \sim 10 \mu\text{K}$)

Deexcitation



Photoassociation spectroscopy

Obtaining spectroscopic signal



Photoassociation spectroscopy

Obtaining spectroscopic signal

“one-color PA”

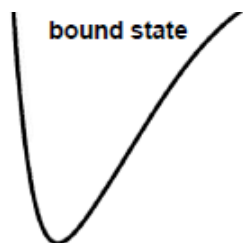
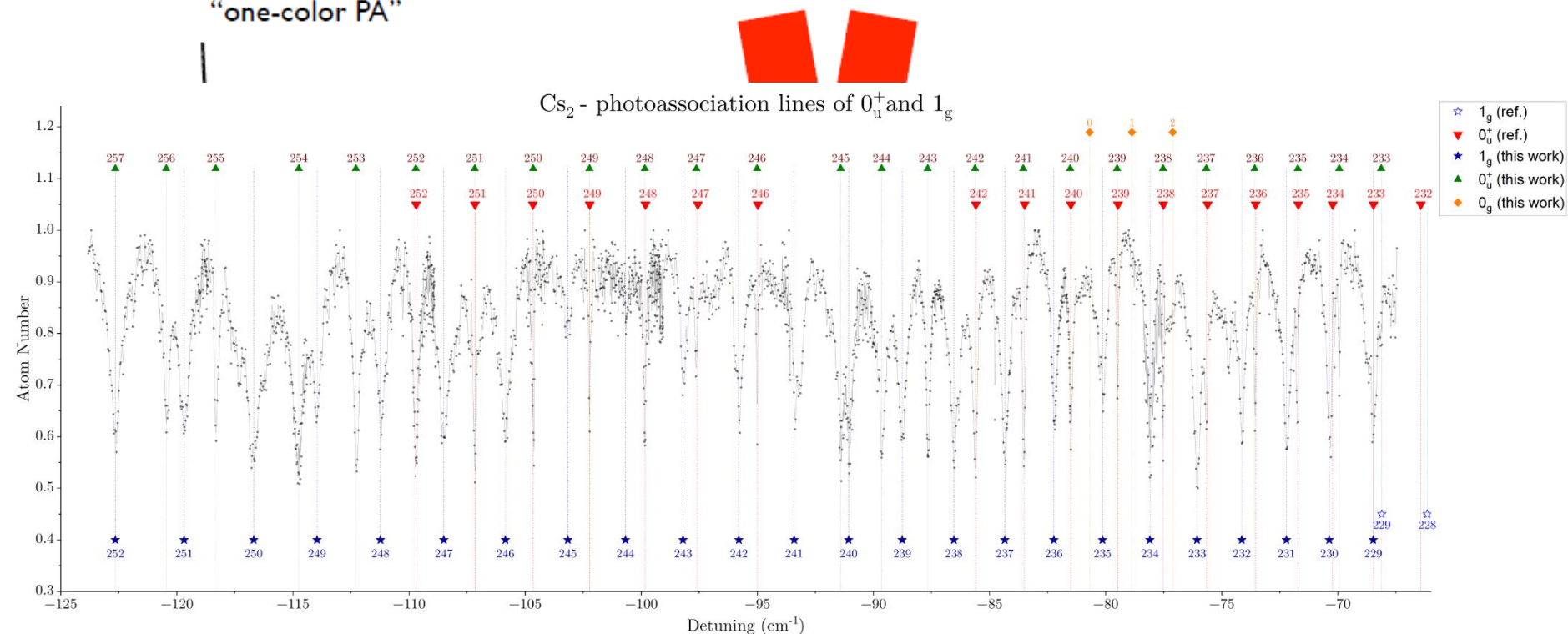


photo-association light

Photoassociation spectroscopy

Obtaining spectroscopic signal

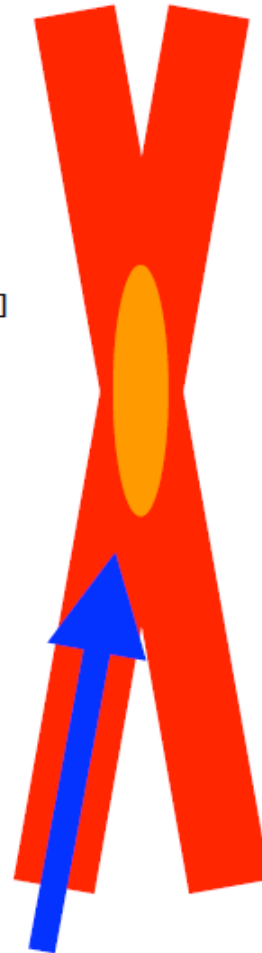
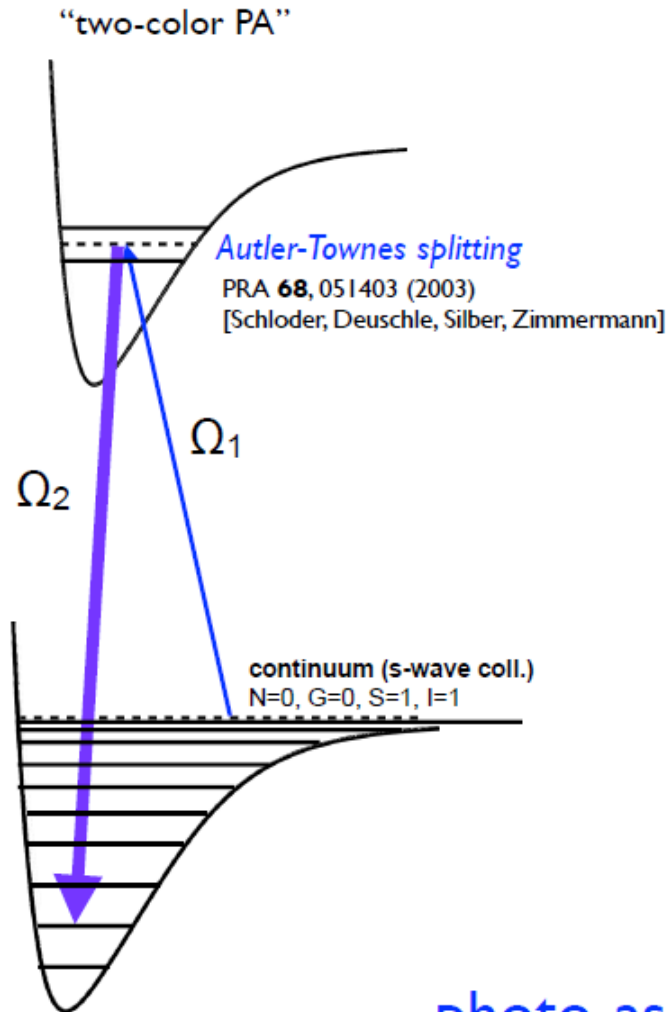
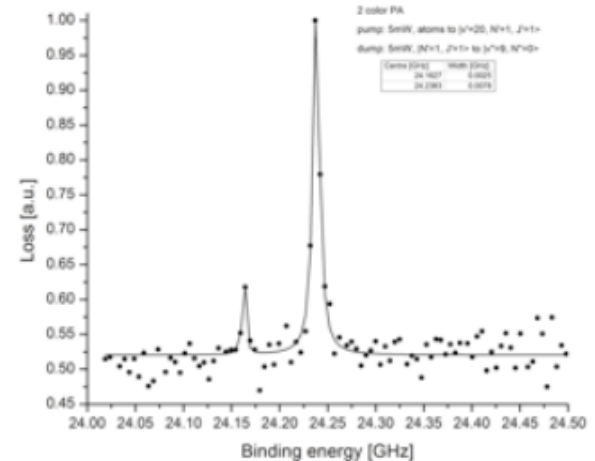


photo-association light

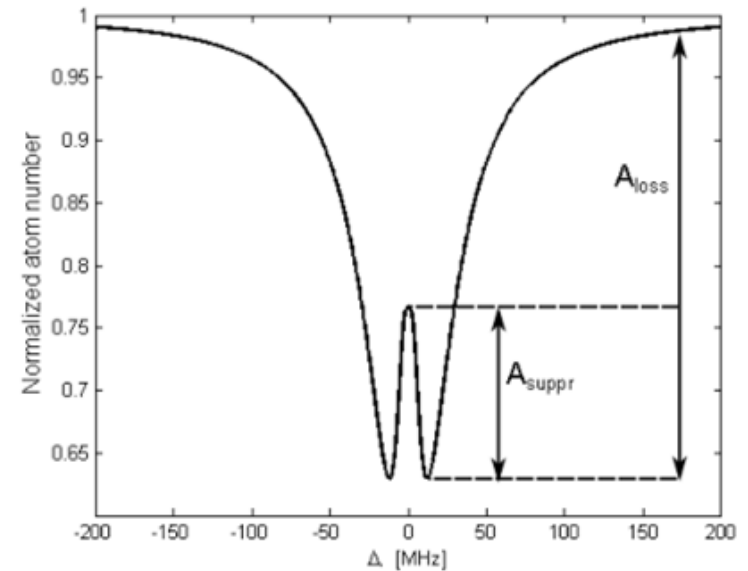
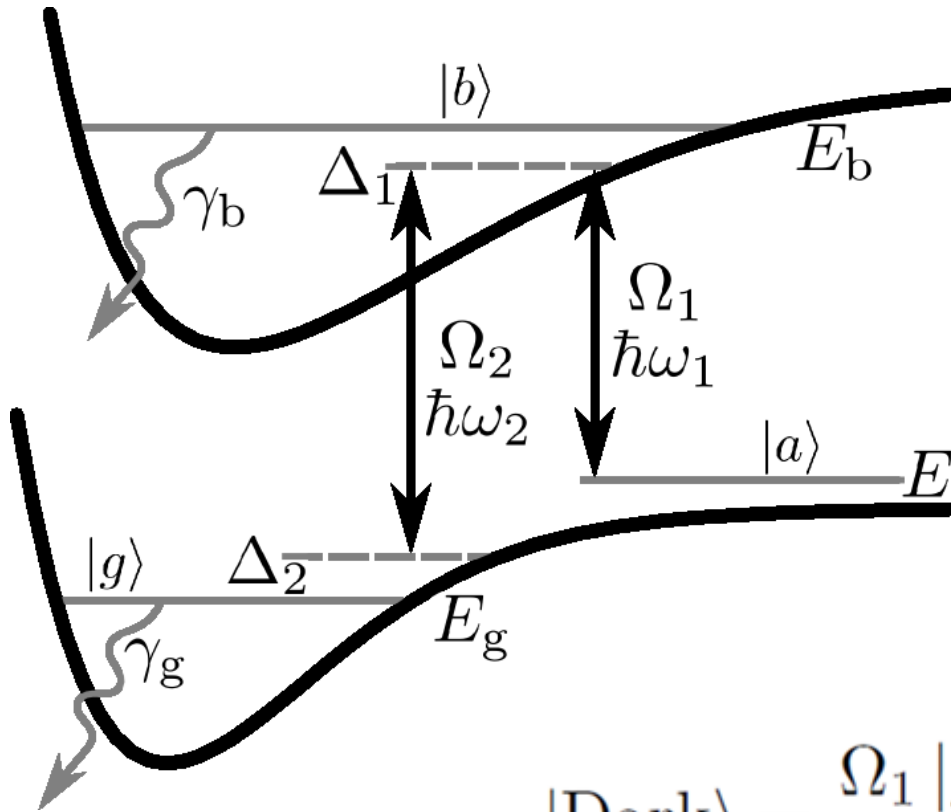
atom loss suppressed



$\Omega_2 - \Omega_1$

width set by coupling induced by Ω_2

Photoassociation spectroscopy



$$|\text{Dark}\rangle = \frac{\Omega_1 |g\rangle - \Omega_2 |a\rangle}{\sqrt{\Omega_1^2 + \Omega_2^2}}$$

Photoassociation spectroscopy

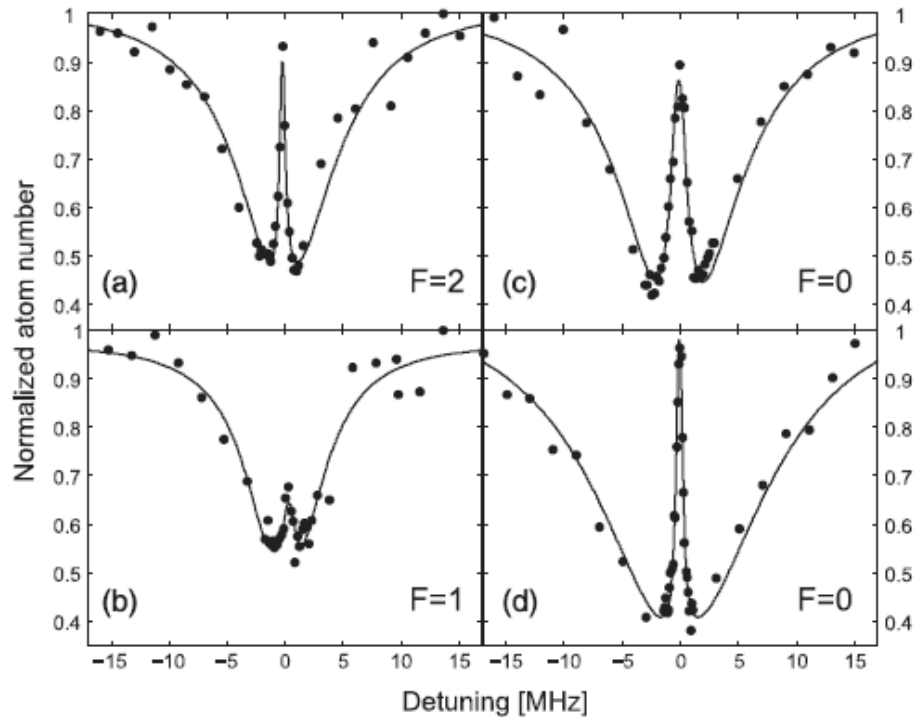
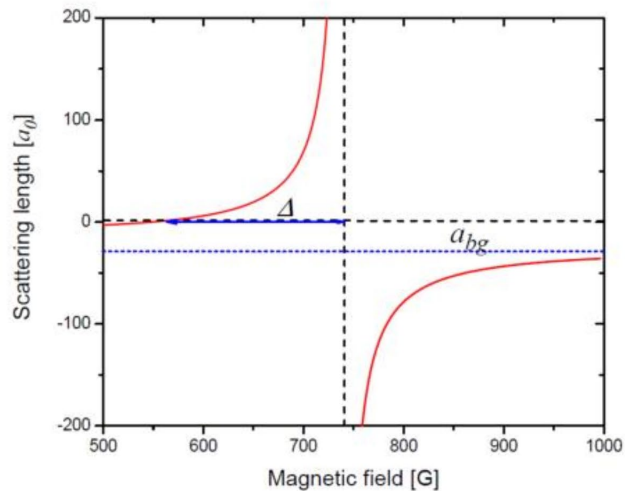
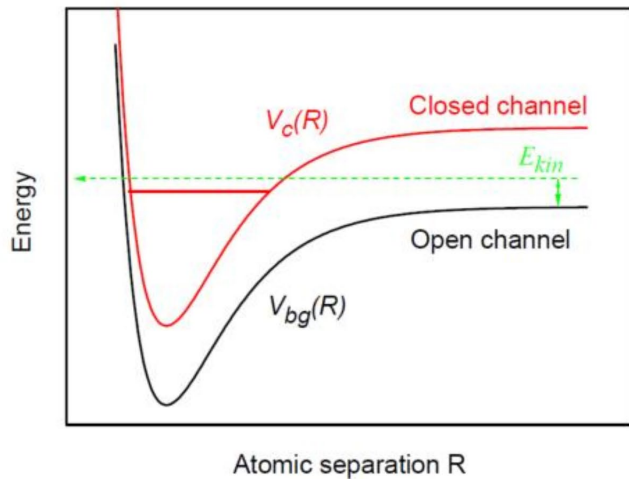


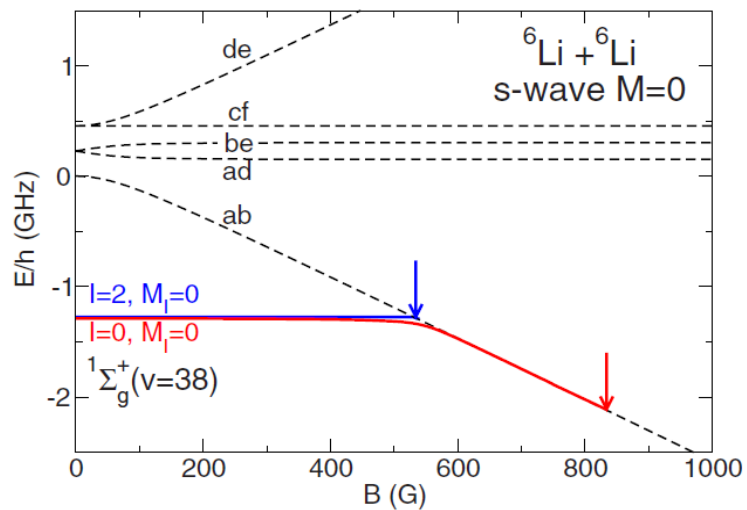
FIG. 2. Dark states in a Fermi gas at 0 G. The atom number is plotted versus the detuning of the probe frequency from the two-photon resonance, where the photon energy difference $h(\nu_2 - \nu_1)$ equals the difference between the energy of the initial free atom state and a bound molecular level and the atom number exhibits a maximum revival. (a)–(c) Corresponding different molecular hyperfine levels ($F'' = 2, 1, 0$) of the $v'' = 9$, $N'' = 0$ level in the $a(1^3\Sigma_u^+)$ potential. We were not able to find parameters that would improve the revival of $F'' = 1$ to above 50%. Spectrum (d) corresponds to the $v'' = 38$, $N'' = I'' = F'' = 0$ level of the $X(1^1\Sigma_g^+)$ potential.

Feshbach spectroscopy

Feshbach resonance

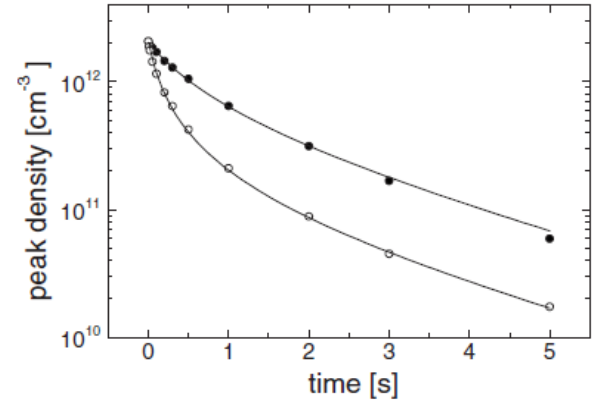
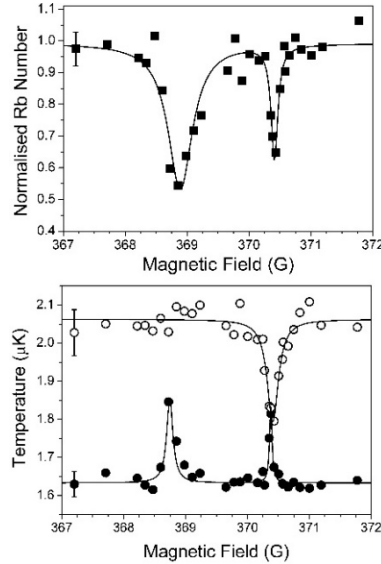
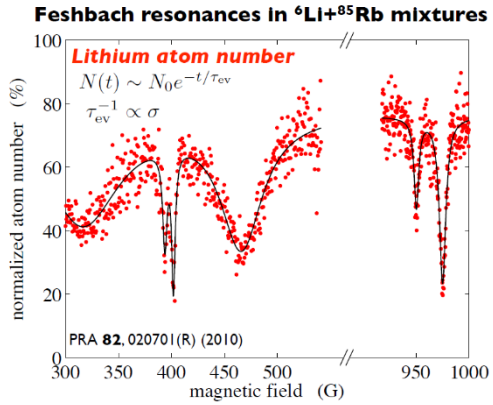


$$a = a_{bg} \left(1 - \frac{\Delta}{B - B_0} \right)$$



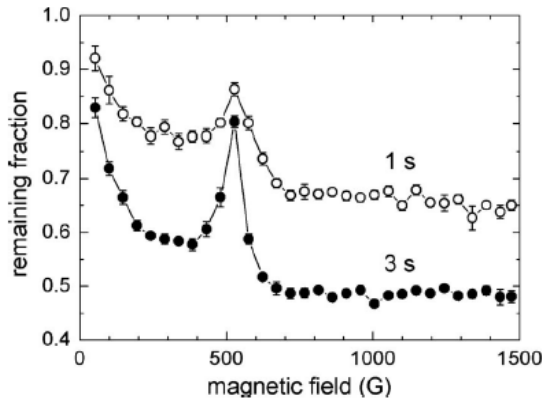
Feshbach spectroscopy

Inelastic loss spectroscopy

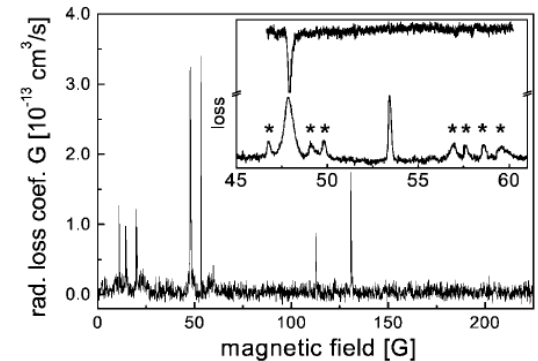
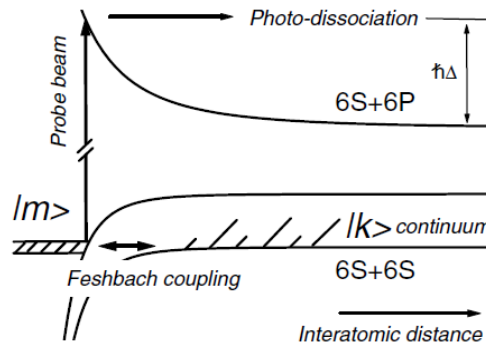


$$\dot{\bar{n}}(t) = -\frac{\bar{n}(t)}{\tau} - L_2 \bar{n}(t)^2 - (4/3)^{3/2} L_3 \bar{n}(t)^3$$

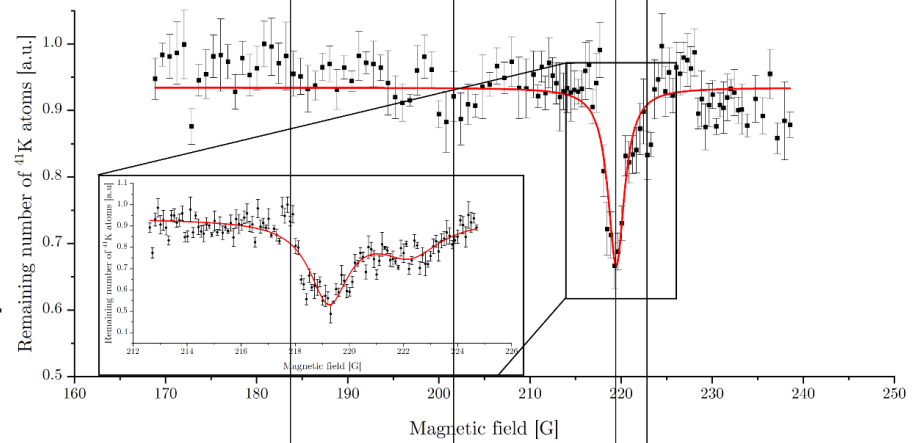
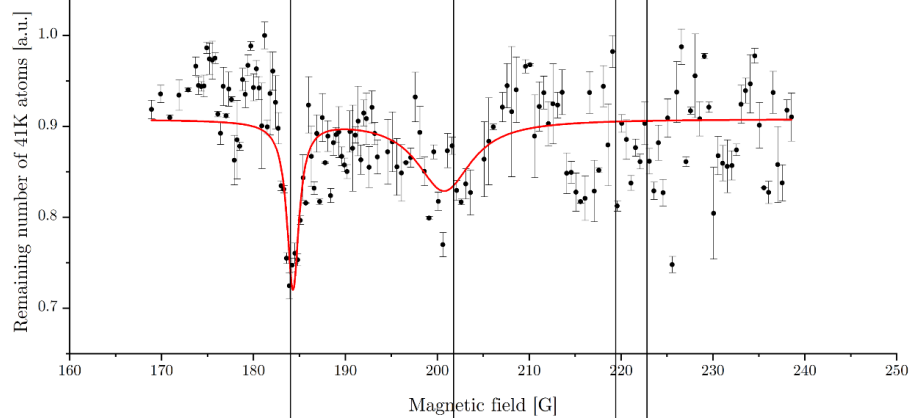
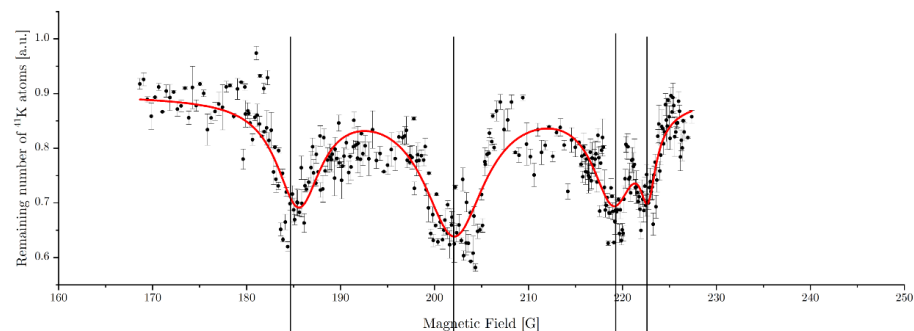
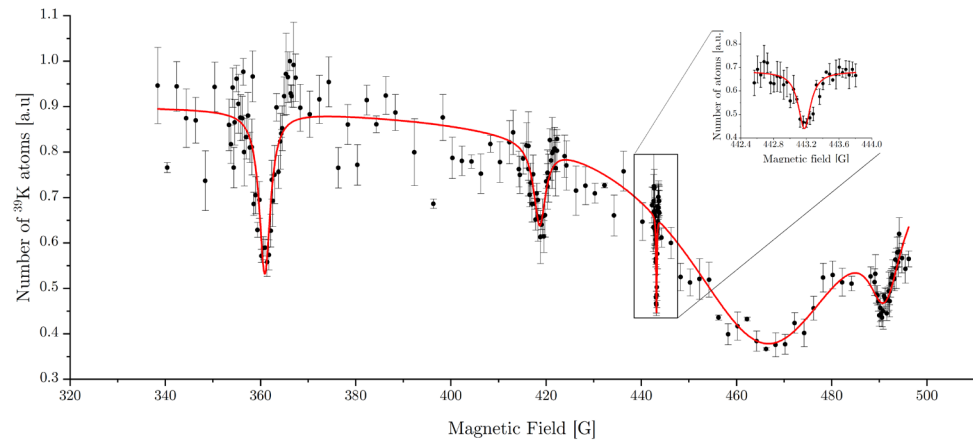
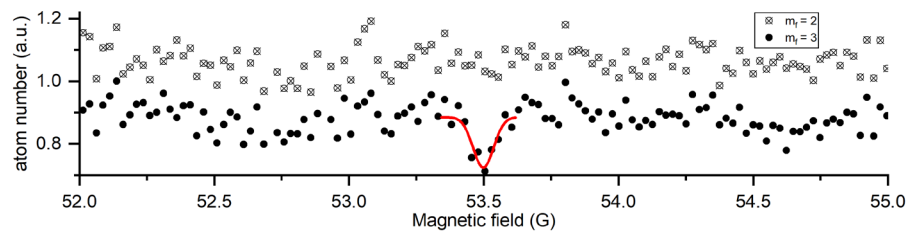
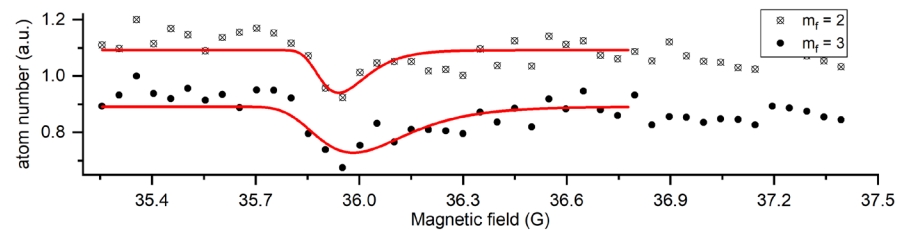
Elastic loss spectroscopy



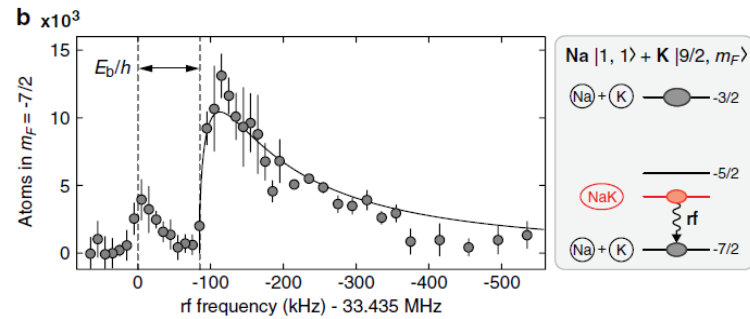
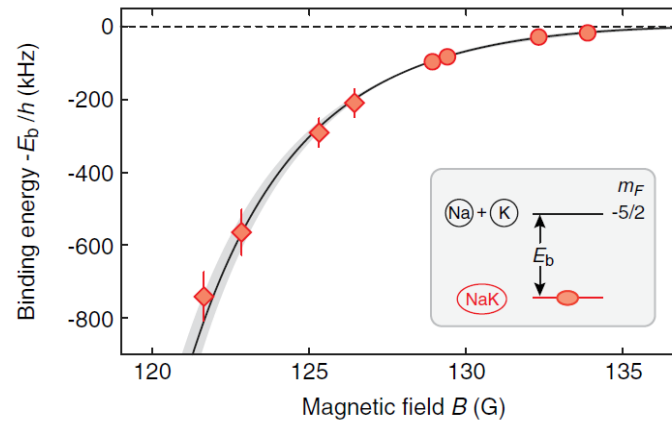
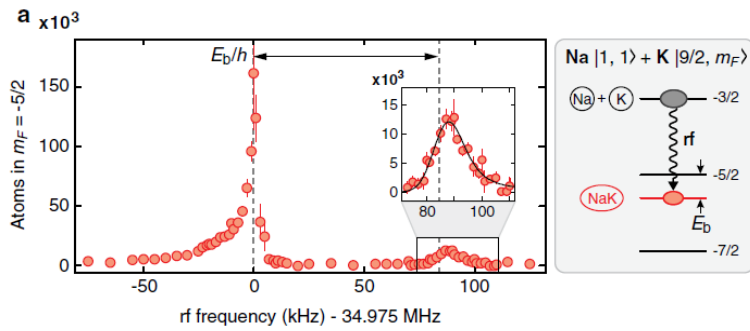
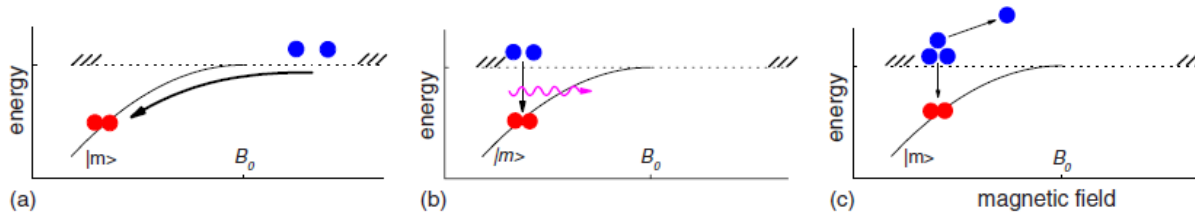
Radiative loss spectroscopy



Feshbach spectroscopy

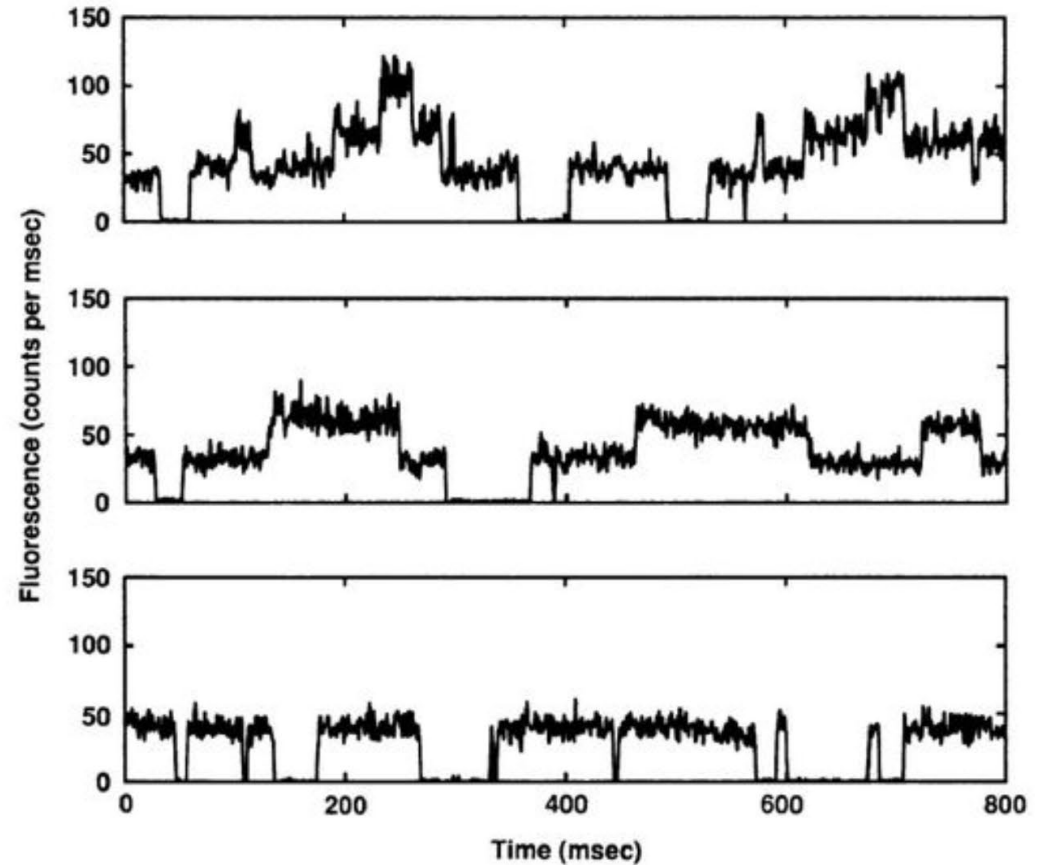
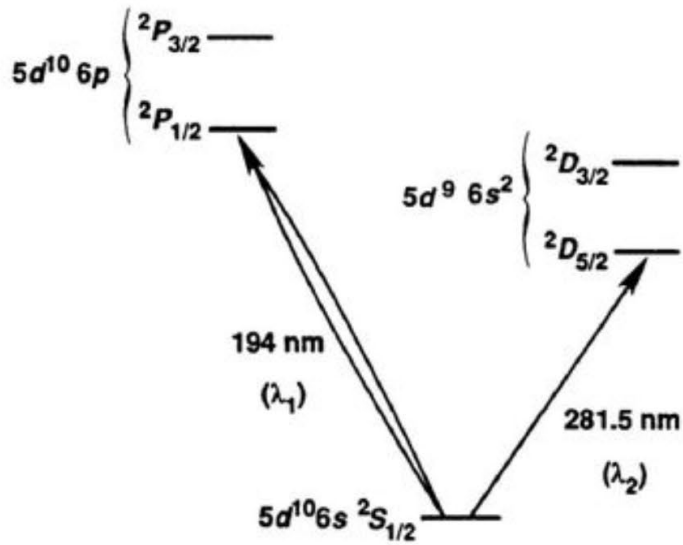


Binding energy measurements



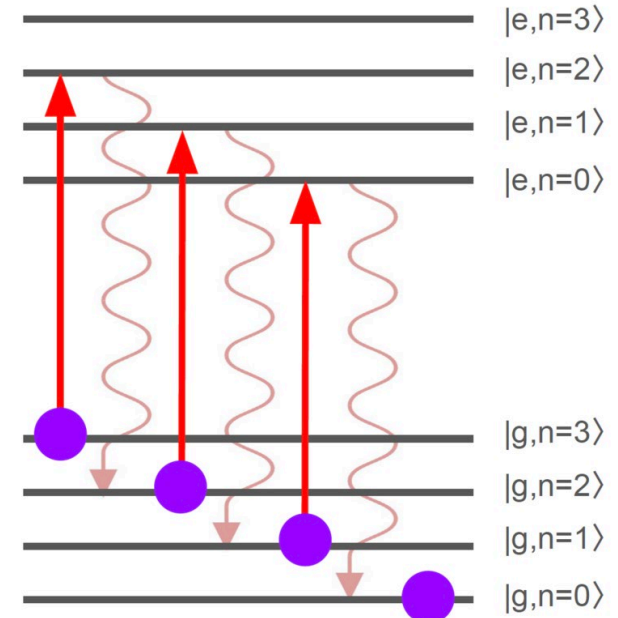
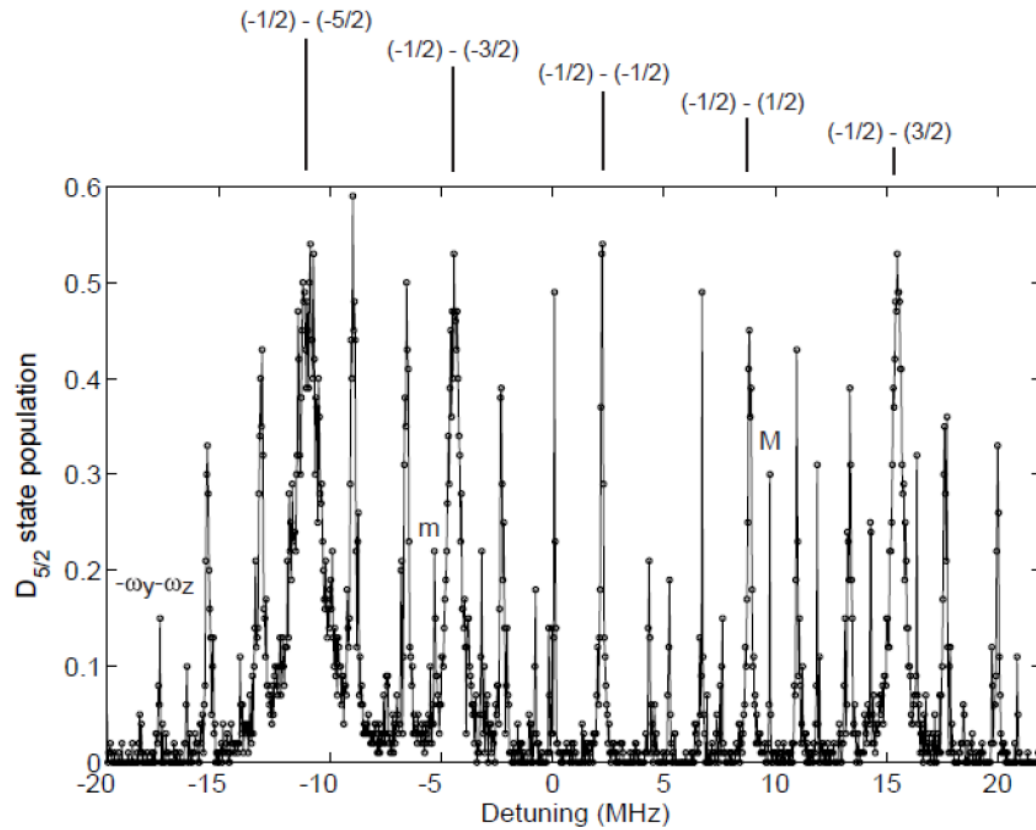
$$E_b \approx \frac{\hbar^2}{2\mu(a-\bar{a})^2}$$

Quantum jump spectroscopy

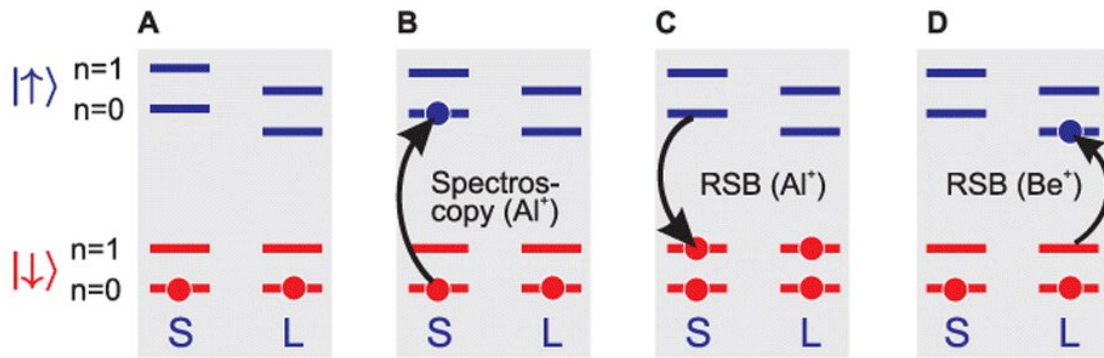


W.M. Itano, J.C. Bergquist, and D.J. Wineland, Science 37, 612 (1987)

Quantum logic spectroscopy



Quantum logic spectroscopy



$$\Psi_0 = |\downarrow\rangle_S |\downarrow\rangle_L |0\rangle_m$$

excite

$$\Psi_0 \rightarrow \Psi_1 = (\alpha|\downarrow\rangle_S + \beta|\uparrow\rangle_S) |\downarrow\rangle_L |0\rangle_m$$

$$= (\alpha|\downarrow\rangle_S |0\rangle_m + \beta|\uparrow\rangle_S |0\rangle_m) |\downarrow\rangle_L$$

red sideband π pulse

$$\Psi_1 \rightarrow \Psi_2 = (\alpha|\downarrow\rangle_S |0\rangle_m + \beta|\downarrow\rangle_S |1\rangle_m) |\downarrow\rangle_L$$

$$= |\downarrow\rangle_S |\downarrow\rangle_L (\alpha|0\rangle_m + \beta|1\rangle_m)$$

red sideband π pulse

$$\Psi_2 \rightarrow \Psi_{\text{final}} = |\downarrow\rangle_S (\alpha|\downarrow\rangle_L + \beta|\uparrow\rangle_L) |0\rangle_m$$

